# Fully nonlinear curvature flow of axially symmetric surfaces with boundary conditions

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Fatemah Mofarreh Fully nonlinear curvature flow of axially symmetric hypersurfaces

# Talk outline



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The setting Previous work Our speeds

Axially symmetric initial hypersurface  $X : [0, a] \times \mathbb{S}^{n-1} \to \mathbb{R}^{n+1}$ ,  $X = (x, u(x)\omega)$ ,

 $u: [0, a] \rightarrow \mathbb{R}^+$  axial graph function. Curvatures

$$\kappa_1 = \frac{-u_{xx}}{\left(1+u_x^2\right)^{\frac{3}{2}}} = \frac{-\left(\arctan u_x\right)_x}{\sqrt{1+u_x^2}}, \ \kappa_2 = \ldots = \kappa_n = \frac{1}{u\sqrt{1+u_x^2}}.$$

Normal curvature flow

$$\left(\frac{\partial X}{\partial t}\left(x,t\right)\right)^{\perp} = -F\left(\mathcal{W}\left(x,t\right)\right)\nu\left(x,t\right),\tag{1}$$

equivalent to graph evolution for  $u : [0, a] \times (0, T)$ 

$$\frac{\partial u}{\partial t} = -\sqrt{1 + u_x^2} F(\mathcal{W}).$$
<sup>(2)</sup>

Example

Speed  $F = H = \kappa_1 + \ldots + \kappa_n$ , mean curvature flow (MCF).

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• (Huisken, '90) MCF, pure Neumann boundary conditions,  $H(\cdot, 0) \ge 0$ . Curvature singularity at  $T < \infty$  is *Type I*:

$$|\mathbf{A}|^2 = \kappa_1^2 + \ldots + \kappa_n^2 \le \frac{C}{T-t}.$$

- (Dzuik-Kawohl, '91) MCF, Neumann/mixed boundary conditions, u(·,0), u<sub>x</sub>(·,0) ≥ 0, H(·,0) ≥ 0. Pinch off at x = 0 at T < ∞.</li>
- (Matioc, '07) MCF, more general boundary conditions:

 $u_{x}(0,t) = 0$  and either  $u_{x}(a,t) = g(t)$  or u(a,t) = h(t)

where *g* and *h* satisfy some natural conditions,

 $u(\cdot,0), u_x(\cdot,0) \ge 0, H(\cdot,0) \ge 0$ . As  $t \to T < \infty$ , either  $\kappa_2^2(0,t) \to \infty$  or  $\kappa_1^2(a,t) \to \infty$ .

 (Escher-Matioc, '10) Periodic MCF, u(·,0), u<sub>x</sub>(·,0) ≥ 0 and H > 0 initially, single point pinch off at T < ∞.</li> Introduction The setting Results Previous wo References Our speeds

Our speed  $F(W) = f(\kappa_1, \ldots, \kappa_n)$  satisfies

i) *f* smooth, symmetric function defined on ℝ<sup>n</sup>.
ii) ∂f/∂κ<sub>i</sub> > 0 for each *i* = 1,..., *n* at every point of Γ.
iii) *f* (*k*κ) = *k f* (κ) for any *k* > 0.
iv) *f* (1,..., 1) = 1.
v) *f* is convex.

For partial singularity classification, need also

$$\lim_{z\to-\infty}f(z,1,\ldots,1)<0.$$

#### Examples

$$f = \left(n + \eta n^{\frac{1}{p}}\right)^{-1} \left[\sum_{i=1}^{n} \kappa_{i} + \eta \left(\sum_{j=1}^{n} \kappa_{j}^{p}\right)^{\frac{1}{p}}\right]$$
for constants  $\eta \in [0, 1)$  and  $p \ge 1$ .

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### Theorem (Short-time existence)

Given  $u_0 \in C^2([0, a])$  compatible with the boundary conditions, there exists  $\delta > 0$  such that there is a unique solution  $u \in C^2([0, a] \times [0, \delta))$  to (2).

### **Remarks:**

- Uniform parabolicity is not required, Condition ii) suffices.
- 2 Higher regularity (smoothing) for a short time, is standard.
- We need C<sup>2</sup> initial data for our partial singularity classification.

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We focus on the Neumann boundary conditions

$$u_{x}(0,t) = 0, u_{x}(a,t) = g(t)$$
 (3)

where g is smooth, non-negative and non-increasing.

#### Lemma

Under the flow (1),

$$u_t(\cdot,0) \leq 0 \text{ that is, } F(\mathcal{W}(\cdot,0)) \geq 0 \implies u_t(\cdot,t) \leq 0.$$

Idea of Proof: Under (2),  $v = u_x$  and  $v = u_t$  satisfy

$$\frac{\partial}{\partial t}v = \frac{\dot{f}^{1}}{1+u_{x}^{2}}v_{xx} - \frac{2\dot{f}^{1}}{\left(1+u_{x}^{2}\right)^{2}}u_{x}u_{xx}v_{x} + \sum_{j=2}^{n}\dot{f}^{j}\frac{v}{u^{2}}.$$

Results follow by the maximum principle, and Hopf Lemma on boundary.

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#### Lemma

 $\begin{array}{ll} \textit{If, in addition to } u_0 > 0, (u_x)_0 \geq 0 \textit{ and } F\left(\mathcal{W}_0\right) \geq 0, \textit{ we have} \\ \arctan g\left(0\right) < \frac{(n-1)a^2}{\int_0^a u_0(x)dx}, & \textit{then } T < \infty. \end{array}$ 

**Idea of Proof:** If instead  $T = \infty$  then u > 0 for all t. Set  $E(t) = \int_0^a u(x, t) dx$ . Then

$$E'(t) = \int_0^a \frac{\partial}{\partial t} u \, dx = -\int_0^a \sqrt{1 + u_x^2} F(\mathcal{W}) \, dx.$$

Since *F* is convex,  $F \geq \frac{1}{n}H$  so

$$E'(t) \leq -\frac{1}{n}\int_0^a \sqrt{1+u_x^2}Hdx \leq \frac{1}{n}\int_0^a (\arctan u_x)_x dx - \frac{n-1}{n}\int_0^a \frac{1}{u}dx.$$

It follows using Hölder's inequality and the assumption that E'(t) < 0, contradicting E(t) > 0 for all t. Hence  $T < \infty$ .

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## Theorem (McCoy, Mofarreh, Williams '13)

Suppose F satisfies given conditions and (2) is accompanied by boundary conditions (3) and  $u_0 > 0$ ,  $(u_0)_x \ge 0$  and  $F \ge 0$ .

If  $\lim_{t\to T} \kappa_1^2(a, T) < \infty$  then  $\lim_{t\to T} \kappa_j^2(0, t) = \infty$  for  $j = 2, \ldots, n$ .

## Idea of Proof:

- Suppose instead lim<sub>t→T</sub> κ<sub>1</sub><sup>2</sup> (a, T) < ∞ and lim<sub>t→T</sub> κ<sub>i</sub><sup>2</sup> (0, T) < ∞.</li>
- We show solution could be extended beyond t = T.
- So we need to show that  $\kappa_1^2$  and  $\kappa_j^2$  remain bounded up to t = T, so  $u(\cdot, T) \in C^2([0, a])$  could be used as new  $u_0$ .
- By assumption  $\lim_{t\to T} u(0,t) = \delta > 0$  so

$$\kappa_{j}^{2}(x,t) = u(x,t)^{-2} \left[ 1 + u_{x}(x,t)^{-2} \right] \le u(x,t)^{-2} \le \delta^{-2}.$$



To bound  $\kappa_1^2$  we need to bound  $u_{xx}$  from above and from below.

## *u<sub>xx</sub>* bound above

By homogeneity, since  $\kappa_2 > 0$ ,

$$f(\kappa_1,\kappa_2,\ldots,\kappa_2) = \kappa_2 f\left(\frac{\kappa_1}{\kappa_2},1,\ldots\right) \ge 0$$

This implies

$$f(z,1,\ldots,1)\geq 0$$

for  $z = \frac{\kappa_1}{\kappa_2}$ . Condition on *f* then implies

$$z = \frac{\kappa_1}{\kappa_2} = \frac{-u \, u_{xx}}{1 + u_x^2} \ge -c_0 \implies \frac{u_{xx}}{1 + u_x^2} \le \frac{c_0}{u} \le \frac{c_0}{\delta}.$$

 $u_{xx}$  bound below The function  $w = -u_t e^{-\lambda t}$  satisfies

$$\frac{\partial}{\partial t}w = \frac{\dot{t}^1}{1+u_x^2}w_{xx} - \frac{2\dot{t}^1}{\left(1+u_x^2\right)^2}u_xu_{xx}w_x + \left(\sum_{j=2}^n\frac{\dot{t}^j}{u^2}-\lambda\right)w.$$

Using the maximum principle, convexity of *F* and the boundary conditions we find that for  $\lambda > \delta^{-2}$ ,

$$w = -u_t e^{-\lambda t} = \sqrt{1 + u_x^2} F e^{-\lambda t} \le \overline{C} (M_0, \delta, T)$$

and so

$$\frac{1}{n}\sqrt{1+u_x^2}\,H\leq\sqrt{1+u_x^2}\,F\leq\overline{C}e^{\lambda T}.$$

Therefore

$$\frac{-u_{xx}}{1+u_x^2}+\frac{n-1}{u}\leq n\,\overline{C}\,e^{\lambda T}$$

providing the lower bound on  $u_{xx}$ . So we have  $\kappa_1^2 \leq C(M_0, \delta, T)$  as required.

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Now consider, for  $0 < k \le 1$  constant, the flow

$$\frac{\partial u}{\partial t} = -\sqrt{1+u_x^2} \, F^k,$$

where F satisfies the same conditions.

Theorem (McCoy, Mofarreh, Williams '13) In the case of pure Neumann boundary conditions and  $u_0 > 0, (u_0)_x \ge 0$  and F > 0, if  $\lim_{t \to T} \kappa_1^2(a, T) < \infty$  then  $\lim_{t \to T} \kappa_j^2(0, t) = \infty$  for j = 2, ..., n.

Remarks on pure Neumann boundary conditions:

- Allow comparison with cylinders to show that  $T < \infty$ .
- Allow application of the maximum principle to the periodic solution to obtain *F* > 0 holds under the flow (not just *F* ≥ 0); needed in the proof for 0 < *k* < 1.</li>



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