Traveling wave dynamics for a one-dimensional constrained Allen-Cahn equation

Goro Akagi

(Mathematical Institute, Tohoku University)

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C. Kuehn (München) and K.-I. Nakamura (Kanazawa)

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1. PDEs with non-decreasing constraints

Non-decreasing constraints

In this talk, "non-decreasing constraints on evolution" mean

$$u(x,t)\geq u(x,s) \hspace{0.1 cm} ext{if} \hspace{0.1 cm} t\geq s,$$

or equivalently,

$$\partial_t u(x,t) \geq 0.$$

Background: Irreversible Phase-field models (e.g., for brittle fracture) Let u = u(x,t) denote a displacement field and z = z(x,t) a phase field, which intuitively means

$$z(x,t) = egin{cases} 1 & ext{if the material is cracked at } (x,t), \ 0 & ext{if the material is not cracked at } (x,t). \end{cases}$$

Then $t \mapsto z(x,t)$ is supposed to be non-decreasing, unlike $t \mapsto u(x,t)$.

Phase-field model for brittle fracture

The phase-field model reads,

$$egin{aligned} 0 &= ext{div} \Big(lpha_arepsilon(z)
abla u \Big) & ext{ in } \Omega imes \mathbb{R}_+, \ \partial_t z &= \Big(arepsilon \Delta z - rac{z}{arepsilon} - lpha_arepsilon'(z) |
abla u|^2 \Big)_+ & ext{ in } \Omega imes \mathbb{R}_+, \end{aligned}$$

which is an irreversible quasi-static evolution, i.e.,

$$0=\partial_u\mathcal{F}_arepsilon(u,z), \hspace{1em} \partial_t z=\Big(-\partial_z\mathcal{F}_arepsilon(u,z)\Big)_+,$$

of the free energy $\mathcal{F}_{\varepsilon}(u, z)$ (= Ambrosio-Tortorelli regularization of the Francfort-Marigo energy) given by

$$\mathcal{F}_arepsilon(u,z) = rac{1}{2}\int_\Omega lpha_arepsilon(z)|
abla u|^2\,\mathrm{d}x + \int_\Omega \left(rac{arepsilon}{2}|
abla z|^2 + rac{1}{2arepsilon}z^2
ight)\mathrm{d}x.$$

[Frémond-Nedjar '96], [Bonetti-Schimperna '04], [Mielke-Roubíček '08], [Knees-Rossi-Zanini '13-], [Takaishi-Kimura '09,'11],...,[A-Schimperna '21]

Gradient flows with non-decreasing constraints

Let us focus on the gradient flow structure with non-decreasing constraints, roughly speaking,

$$\partial_t u(t) = \Big(- \partial J(u(t)) \Big)_+ \geq 0$$

for some (possibly non-convex) functional $J: X \to \mathbb{R}$, say $X = L^2(\Omega)$. Equivalently, it can be rewritten as

$$\partial_t u(t) = -\partial J(u(t)) - \mu,$$

where μ can be characterized as

$$\mu\in\partial I_{[0,+\infty)}(\partial_t u)$$
 and $\mu=-\Big(-\partial J(u(t))\Big)_-.$

Stabilization \leftrightarrow **Constraint**

Aim of this talk

Address ourselves onto <u>simpler PDEs</u> with non-decreasing constraints and discuss asymptotic behavior of solutions.

In particular, we shall discuss traveling wave dynamics for the 1D Allen-Cahn equation with non-decreasing constraints,

$$(st) \qquad \qquad u_t = \Big(u_{xx} - W'(u) \Big)_+ \; ext{ in } \mathbb{R} imes \mathbb{R}_+.$$

cf.) [A-Efendiev '19] Cauchy-Dirichlet problem in bounded domains of \mathbb{R}^N

- well-posedness, (partial) smoothing effect
- energy-dissipation estimates, absorbing set, global attractor
- reformulation of (*) as an obstacle problem
- convergence to equilibria as $t
 ightarrow +\infty$ and steady-state equation

2. Traveling wave dynamics

Allen-Cahn equation in \mathbb{R}

Let us recall the Cauchy problem for the classical Allen-Cahn equation,

$$(\mathsf{AC})egin{cases} u_t = u_{xx} - f(u) & ext{ in } \mathbb{R} imes \mathbb{R}_+, \ u|_{t=0} = u_0 & ext{ in } \mathbb{R}, \end{cases}$$

where f satisfies

$$egin{cases} f(a_\pm) = f(a_0) = 0, & f'(a_\pm) > 0, \ f > 0 & ext{in} \; (a_-, a_0), & f < 0 & ext{in} \; (a_0, a_+), \end{cases}$$

for some $a_- < a_0 < a_+$, i.e., f = W' with a double-well potential W.

Allen-Cahn equation in \mathbb{R}

Let us recall the Cauchy problem for the classical Allen-Cahn equation,

$$(\mathsf{AC}) egin{cases} u_t = u_{xx} - f(u) & ext{ in } \mathbb{R} imes \mathbb{R}_+, \ u_{t=0} = u_0 & ext{ in } \mathbb{R}. \end{cases}$$

phase separation model (e.g., binary alloy),

 $\spadesuit L^2$ gradient flow of the free energy

$$J(u):=rac{1}{2}\int_{\mathbb{R}}|\partial_{x}u(x)|^{2}\,\mathrm{d}x+\int_{\mathbb{R}}W(u(x))\,\mathrm{d}x$$

where W' = f. Namely,

$$(\mathsf{AC}) \Leftrightarrow u_t = -\partial J(u) ext{ in } L^2(\mathbb{R}), ext{ } t > 0.$$

Hence $t \mapsto J(u(t))$ is non-increasing.

Traveling wave solutions for (AC)

The traveling wave solution $u(x,t) = \phi(x - ct)$ is characterized by a profile function ϕ and a velocity constant c satisfying

$$egin{cases} -c\phi'=\phi''-f(\phi)& ext{in }\mathbb{R},\ \phi(\xi) o a_{\pm}& ext{as }x o\pm\infty \end{cases}$$

and

$$c egin{cases} < 0 & ext{if} \ W(a_+) < W(a_-), \ = 0 & ext{if} \ W(a_-) = W(a_+), \ > 0 & ext{if} \ W(a_-) < W(a_+). \end{cases}$$



Traveling wave solutions for (AC)

[Fife-McLeod '72]

• Existence and "uniqueness" of TWs

Phase-plain analysis

• Exponential stability of TWs: if

$$u_0(x)\in [a_-,a_+], \hspace{1em} \limsup_{x
ightarrow -\infty} u_0(x) < a_0, \hspace{1em} \liminf_{x
ightarrow +\infty} u_0(x) > a_0,$$

then there exist constants $x_0 \in \mathbb{R}$, $K, \kappa > 0$ such that

$$\|u(\cdot,t)-\phi(\cdot-ct-x_0)\|_{L^{\infty}(\mathbb{R})}\leq K\mathrm{e}^{-\kappa t} ext{ for all } t\geq 0.$$

Schauder estimate, precompactness of $\{u(\cdot+ct,t)\colon t\geq 0\}$, sub / supersolution method

[X. Chen '92] Non-local reaction, Sub / supersolution method only

Constrained Allen-Cahn equation

Now, we shall consider

$$\mathsf{(AC)}_+ egin{cases} u_t = ig(u_{xx} - f(u) ig)_+ & ext{ in } \mathbb{R} imes \mathbb{R}_+, \ u_{|_{t=0}} = u_0 & ext{ in } \mathbb{R}, \end{cases}$$

where f satisfies

$$egin{cases} f(a_\pm) = f(a_0) = 0, & f'(a_\pm) > 0, \ f > 0 & ext{in} \ (a_-, a_0), & f < 0 & ext{in} \ (a_0, a_+) \end{cases}$$

for some $a_- < a_0 < a_+$ and

$$(s)_+ := \max\{s, 0\} \ge 0.$$

Constrained Allen-Cahn equation

Now, we shall consider

$$(\mathsf{AC})_+ egin{cases} u_t = ig(u_{xx} - f(u) ig)_+ & ext{in } \mathbb{R} imes \mathbb{R}_+, \ = - \partial J(u) \ u|_{t=0} = u_0 & ext{in } \mathbb{R}, \end{cases}$$

Then $t \mapsto J(u(t))$ is still non-increasing.

Moreover, $(AC)_+$ is equivalent to

$$ullet \ u_t=-\,\partial J(u)-\mu, \quad \mu\in\partial I_{[0,+\infty)}(u_t), \quad u|_{t=0}=u_0,$$

$$egin{aligned} & egin{aligned} & egin{aligned} & egin{aligned} & egin{aligned} & eta &$$

Traveling wave dynamics for $(AC)_+$

We shall discuss

- existence and uniqueness of traveling wave solutions,
- convergence to a traveling wave solution.

In what follows, we may restrict ourselves to a balanced potential:

$$a_{\pm}=\pm 1, \ \ a_0=0, \ \ W(u)=rac{1}{4}(u^2-1)^2,$$

for which the TW of (AC) fulfills

c=0 and $\phi(\pm\infty)=\pm 1.$

3. Existence and uniqueness of traveling waves

Balanced double-well potential



Balanced double-well potential W(u)

Heuristic construction

Substitute $u(x,t) = \phi(x-ct)$ to $(\mathsf{AC})_+$. Then a profile equation reads,

$$-c\phi'=\Bigl(\phi''-f(\phi)\Bigr)_+.$$

How can we solve it ?

Instead, we shall derive an alternative profile equations...

Heuristic construction

Suppose that $u(-\infty) \equiv \alpha \in (-1, 0)$, which is a steady-state for $(AC)_+$. (i) Due to the non-decreasing constraint $u_t \ge 0$,

 $u(x,t)\geq lpha.$

Then the region $[u < \alpha]$ on the graph of W(u) is prohibited.

(ii) Analogously to the reformulation, u may solve the obstacle problem,

$$u_t - u_{xx} + \underbrace{f(u)}_{=W'(u)} + \partial I_{[lpha,+\infty)}(u)
i 0 ext{ in } \mathbb{R} imes \mathbb{R}.$$

Roughly speaking,

$$W'(u) + \partial I_{[lpha,+\infty)}(u) = \partial ig(W + \underbrace{I_{[lpha,+\infty)}}_{ ext{truncation }!}ig)(u).$$

Heuristic construction



Truncated balanced double-well potential W(u)

Profile equation for traveling waves

Let $lpha \in (-1,0)$ and substitute $u = \phi(x-ct)$. Then

$$(*) \qquad -c\phi' - \phi'' + \underbrace{W'(\phi) + \partial I_{[\alpha,\infty)}(\phi)}_{= \partial(W + I_{[\alpha,\infty)})(\phi)} \ni 0 \text{ in } \mathbb{R}$$

is derived along with

$$\phi(-\infty) = \alpha$$
 and $\phi(+\infty) = 1$ and $\phi' \ge 0$

as a profile equation.

Indeed, let ϕ_{α} be a solution of (*) for some $c = c_{\alpha}$.

Then $u(x,t) = \phi_{lpha}(x-c_{lpha}t)$ solves $(\mathsf{AC})_+$

and $c_{\alpha} < 0$ if $\alpha \neq -1$.



- Theorem 1 (Traveling wave solutions [A-K-N]) For each $\alpha \in (-1,0)$, $(AC)_+$ has a solution $u(x,t) = \phi_{\alpha}(x - c_{\alpha}t)$ for some profile function $\phi_{\alpha}(\xi)$ and a velocity constant c_{α} satisfying

$$\lim_{\xi o+\infty}\phi_lpha(\xi)=1,\quad \lim_{\xi o-\infty}\phi_lpha(\xi)=lpha.$$

Moreover, it holds that

(i)
$$-c\phi'_{\alpha} - \phi''_{\alpha} + f(\phi_{\alpha}) + \partial I_{[\alpha,\infty)}(\phi_{\alpha}) \ni 0$$
 in \mathbb{R} ,
 $\Leftrightarrow \min \{\phi_{\alpha} - \alpha, -c_{\alpha}\phi'_{\alpha} - \phi''_{\alpha} + f(\phi_{\alpha})\} = 0$ in \mathbb{R} ,
(ii) $\phi_{\alpha} \in W^{2,\infty}(\mathbb{R}), \quad \phi'_{\alpha} \in L^{2}(\mathbb{R}),$
(iii) $\alpha \leq \phi_{\alpha} < 1, \quad 0 \leq \phi'_{\alpha} < +\infty$ in $\mathbb{R}, \quad -\infty < c_{\alpha} < 0$,

 \sim Theorem 1 (Traveling wave solutions [A-K-N] (contd.)) — There exists $s_0 \in \mathbb{R}$ such that

 $\phi_lpha(s)=lpha \ \ ext{on} \ (-\infty,s_0], \ \ \ \phi_lpha(s)>lpha \ \ \ ext{on} \ (s_0,\infty).$

Furthermore, $-c_{lpha}\phi_{lpha}' - \phi_{lpha}'' + f(\phi_{lpha}) = 0$ in $(s_0, +\infty)$, and hence,

$$\phi_lpha''(s_0-0)=0$$
 and $\phi_lpha''(s_0+0)=f(lpha)>0,$

which implies $\phi \not\in C^2(\mathbb{R})$.

In what follows, we set $s_0 = 0$ by translation. Hence

$$\phi_lpha=lpha \;\; ext{on}\; (-\infty,0], \;\;\; \phi_lpha>lpha \;\; ext{on}\; (0,\infty).$$



Uniqueness of profile and velocity

- Theorem 2 (Uniqueness of traveling waves [A-K-N]) — Concerning traveling wave solutions discussed in Theorem 1, it holds that (i) the velocity c_{α} is unique for each α ,

(ii) the profile function ϕ_{α} is unique (up to translation) for each α .



4. Convergence to traveling waves

Question

Question

Can we also prove the (exponential) convergence of solutions for $(AC)_+$ to traveling waves as in the classical Allen-Cahn equation ?

We cannot expect stability of traveling wave solutions.

- Proposition 3 (Instability of traveling waves [A-K-N]) — For each $\alpha \in (a_-, a_0)$, the traveling wave solution $\phi_{\alpha}(x - c_{\alpha}t)$ of $(AC)_+$ is unstable in $L^{\infty}(\mathbb{R})$.

Hence the basin of attraction of the traveling wave solution for $(AC)_+$ is smaller than those for (AC).

Hypotheses for initial data

For $\alpha \in (a_-, a_0)$, we assume (with $a_{\pm} = \pm 1$ and $a_0 = 0$): $(\mathsf{H})_{\alpha} \begin{cases} u_0 \in H^2_{\mathrm{loc}}(\mathbb{R}), & \liminf_{x \to +\infty} u_0(x) > a_0, \\ a_- \leq \inf_{x \in \mathbb{R}} u_0(x) \leq \sup_{x \in \mathbb{R}} u_0(x) < a_+, \\ u_0 \equiv \alpha \quad \mathrm{on} \ (-\infty, \xi_1] \quad \mathrm{for \ some} \ \xi_1 \in \mathbb{R}. \end{cases}$



Then one can define

 $r(t):=\sup\{r\in \mathbb{R}\colon u(x,t)=lpha ext{ for all } x\leq r\}\in \mathbb{R} ext{ for } t\geq 0.$

Main result

Theorem 4 (Exponential convergence to a TW [A-K-N]) — Let $\alpha \in (a_-, a_0)$ be such that $f'(\alpha) > 0$ and assume that u_0 satisfies $(\mathsf{H})_{\alpha}$. Let u = u(x, t) be the L^2_{loc} solution to $(\mathsf{AC})_+$ for the initial datum u_0 . Then there exist $x_0 \in \mathbb{R}$ and $K, \kappa > 0$ such that $\|u(\cdot, t) - \phi_{\alpha}(\cdot - c_{\alpha}t - x_0)\|_{L^{\infty}(\mathbb{R})} \leq Ke^{-\kappa t}$ for all $t \geq 0$. Set $r(t) := \sup\{r > 0 : u(\cdot, t) \equiv \alpha \text{ on } (-\infty, r]\}$. Then

$$|r(t)-c_lpha t-x_0|\lesssim {
m e}^{-rac{\kappa}{2}t}$$
 as $t o+\infty.$

Supersolutions of $(AC)_+$ cannot decrease !

We note that

$$U_t \geq \left(U_{xx} - f(U)
ight)_+ \geq 0$$

implies

U(x,t) cannot decrease in time !

Hence the sub- and supersolution method does not work well for $(AC)_+$ (cf. Fife-McLeod & Chen).

 \Rightarrow Our proof consists of "4 phases".

Rough outline of proof

Phases 1 & 2 Reduction to a simplified system

There exists $t_1 > 0$ such that, for all $t \ge t_1$, $u(x,t) = u_0(x)$ if and only if $u(x,t) = \alpha$. Hence $(AC)_+$ is reduced to $(AC)_{\alpha} \min \{u - \alpha, u_t - u_{xx} + f(u)\} = 0$ in \mathbb{R} .

Phase 3Quasi-convergence of the orbit $\mathcal{O} = \{u(\cdot + c_{\alpha}t, t) : t \geq t_1\}$ to the limit $\phi_{\alpha}(\cdot - x_0)$ unif. in \mathbb{R} for some $x_0 \in \mathbb{R}$ Energy methodPhase 4Exponential convergence of $u(\cdot, t)$ to $\phi_{\alpha}(\cdot - c_{\alpha}t - x_0)$ uniformly in \mathbb{R} as $t \to +\infty$ Sub- and supersolution method

Initial Phase

Claim $\exists t_1 > 0 \ ; \ \inf_{x \in \mathbb{R}} u(x,t_1) > a_-.$

Let u_{ac} be the unique solution to (AC) with the same datum u_0 . Moreover,

$$u_t = \Big(u_{xx} - f(u)\Big)_+ \geq u_{xx} - f(u) \hspace{0.2cm} ext{in} \hspace{0.2cm} \mathbb{R} imes \mathbb{R}_+.$$

By comparison principle,

$$u(x,t) \geq u_{ac}(x,t)
ightarrow \phi_{ac}(x) \ \ {
m as} \ \ t
ightarrow +\infty,$$

where ϕ_{ac} is a TW with c = 0 connecting a_{\pm} at $x = \pm \infty$.



Let $u_{ac}(x,t)$ be the solution to (AC) with $u_{ac}(x,0) = u_0(x)$.



Let $u_{ac}(x,t)$ be the solution to (AC) with $u_{ac}(x,0) = u_0(x)$.



Then $u_{ac}(x,t)$ converges to a layer solution $\phi_{ac}(x)$ (with c=0).



The solution u(x,t) of $(\mathsf{AC})_+$ is a supersolution to (AC). Hence $u(x,t) \ge u_{ac}(x,t)$.

Second Phase

- Lemma 5 (Reduction to a constant obstacle problem) -There exists $t_2 > t_1$ such that, for all $t > t_2$,

$$u(x,t)=u_0(x) \quad \Leftrightarrow \quad u(x,t)=lpha.$$

We employ a subsolution to $(AC)_+$ given by

$$U_\gamma(x,t):=\phi_\gamma(x-c_\gamma t-\sigma\delta(1-\mathrm{e}^{-eta t})-h^-)-\delta\mathrm{e}^{-eta t}$$

for some $h^- \in \mathbb{R}$, $eta, \delta, \sigma > 0$ and $\gamma \in (a_-, a_0)$ satisfying

$$a_- < \gamma - \delta < \inf_{x \in \mathbb{R}} u(x,t_1).$$

Then we assure that $c_{\gamma} < 0$.



$$\inf_{x\in \mathbb{R}} u(x,t_1) > a_-$$



Eventually, $u(x, t + s) \ge U_{\beta,\gamma}(x, s) > \alpha = u_0(x)$ for all $x \ge \xi_1$. Therefore $u(x, t) = u_0(x) \Leftrightarrow u(x, t) = \alpha$.

Thanks to the reformulation, we assure that

$$(\mathsf{AC})_+ \Leftrightarrow (\mathsf{AC})_{lpha} \min \{u - lpha, u_t - u_{xx} + f(u)\} = 0$$
 in \mathbb{R} for all $t > t_2$.

Third Phase

– Lemma 6 (Quasi-convergence of $u(\cdot - c_lpha t, t))$

There exist a sequence $t_n o +\infty$ and $\xi \in \mathbb{R}$ such that

$$\|u(\cdot,t_n)-\phi_lpha(\cdot-c_lpha t_n+\xi)\|_{L^\infty(\mathbb{R})} o 0.$$

We prove the precompactness of $\{u(\cdot - c_{\alpha}t, t) : t > t_2\}$ in $H^1_{loc}(\mathbb{R})$ by developing local energy estimates for $(AC)_{\alpha}$ and identify the limit of the orbit.

Thereby, for any $\delta > 0$ (small), one can take $n_{\delta} \in \mathbb{N}$ such that

$$\phi_lpha(x-c_lpha t_{n_\delta})-\delta\leq u(x,t_{n_\delta})\leq \phi_lpha(x-c_lpha t_{n_\delta})+\delta$$

for any $x \in \mathbb{R}$ (after a suitable translation).

Set

$$v(y,t):=u(y+c_lpha t,t)-a_+\in [lpha-a_+,0], \ y\in \mathbb{R}_+, \ t>t_2.$$

Then $(\mathsf{AC})_lpha$ implies

$$\partial_t v - \partial_y^2 v + \eta + f(v+a_+) = c_lpha \partial_y v, \hspace{1em} \eta \in \partial I_{[lpha,\infty)}(v+a_+)$$

in $\mathbb{R}_+ \times \mathbb{R}_+$. We can assume WLOG (by translation) that

$$v(0,t) = lpha - a_+, \ \ \partial_y v(0,t) = 0, \ \ v(y,0) = u(y,0).$$

A key step is to establish local-energy estimates for v(y, t).

Step 1 | Based on [A-Efendiev '19], we can prove that

$$\sup_{t\geq 0}\int_{\mathbb{R}}|\eta(\cdot,t)|^2\rho\,\mathrm{d} x\leq \int_{\mathbb{R}}|(\partial_x^2u_0-f(u_0))_-|^2\rho\,\mathrm{d} x\quad \forall\rho\in C_c^\infty(\mathbb{R}).$$

Step 2 (Caccioppoli type estimate) Let $\zeta_R \in C^\infty_c(\mathbb{R})$ be such that

$$\zeta_R\equiv 1 \hspace{0.2cm} ext{in} \hspace{0.1cm} [0,R], \hspace{0.2cm} \zeta_R\equiv 0 \hspace{0.2cm} ext{in} \hspace{0.1cm} [2R,+\infty), \hspace{0.2cm} \|\zeta_R'\|_{L^\infty(\mathbb{R}_+)}\leq rac{2}{R}.$$

Test $(\mathsf{AC})_{\alpha}$ by $\mathrm{e}^{-\lambda t} \mathrm{e}^{c_{\alpha} y} v \zeta_R^2$. Then

$$\begin{split} \frac{1}{2} \mathrm{e}^{-2\lambda t} \int_{0}^{+\infty} \mathrm{e}^{c_{\alpha} y} v(\cdot, t)^{2} \zeta_{R}^{2} \,\mathrm{d}y + \frac{1}{2} \int_{0}^{t} \mathrm{e}^{-2\lambda \tau} \left(\int_{0}^{+\infty} \mathrm{e}^{c_{\alpha} y} |\partial_{y} v|^{2} \zeta_{R}^{2} \,\mathrm{d}y \right) \mathrm{d}\tau \\ & \leq \frac{1}{2|c_{\alpha}|} \|v(\cdot, 0)\|_{L^{\infty}(\mathbb{R}_{+})}^{2} + \frac{4}{\lambda |c_{\alpha}|R^{2}} \|v\|_{L^{\infty}(Q_{T})}^{2}, \end{split}$$

where we used the fact that $f'' \geq -\lambda$ and $v\eta \geq 0$. Let $R \to +\infty$. Then

$$\int_0^t \mathrm{e}^{-2\lambda au} \left(\int_0^{+\infty} \mathrm{e}^{c_lpha y} |\partial_y v|^2 \,\mathrm{d} y
ight) \mathrm{d} au \leq C.$$

Step 3 (Weighted energy estimate for $\partial_t v$) Test $(AC)_{\alpha}$ by $e^{c_{\alpha} y} \partial_t v$. Then

$$\int_0^\infty e^{c_\alpha y} |\partial_t v|^2 \, \mathrm{d}y + \frac{\mathrm{d}}{\mathrm{d}t} \left(\frac{1}{2} \int_0^\infty e^{c_\alpha y} |\partial_y v|^2 \, \mathrm{d}y + \int_0^\infty e^{c_\alpha y} h(v) \, \mathrm{d}y \right) = 0.$$
$$= E(v(\cdot, t))$$

where $h(v) = \hat{f}(v + a_+) \ge 0$, since $\eta \partial_t v \equiv 0$ by $\eta \in \partial I_{[0,+\infty)}(\partial_t u)$. Integrate it over (τ_0, t) with $\tau_0 > 0$. It then follows from Step 1 that

$$\int_{ au_0}^t \int_0^\infty \mathrm{e}^{c_lpha y} |\partial_t v(y, au)|^2 \,\mathrm{d}y \mathrm{d} au + E(v(\cdot,t)) \leq E(v(\cdot, au_0)) \quad orall t \geq au_0.$$

Step 4 (Quasi-convergence local in space) One can take $t_n \to \infty$ such that

 $\partial_t v(\cdot, t_n) o 0$ strongly in $L^2(\mathbb{R}_+; \mathrm{e}^{c_{\alpha} y} \mathrm{d} y),$

$$egin{aligned} &\eta(\cdot,t_n) o \eta_\infty & ext{weakly in } L^2(0,R), \ &v(\cdot,t_n) o \psi & ext{weakly in } H^2(0,R), \ & ext{ strongly in } C^1([0,R]) \end{aligned}$$

for some $\psi \in H^2_{\mathrm{loc}}(\mathbb{R}_+)$ and any R>0. Set

$$\phi(y) = egin{cases} \psi(y) + a_+ & ext{if} \ y \geq 0, \ lpha & ext{if} \ y < 0, \end{cases}$$

which then solves

$$-\phi''+f(\phi)+I_{[lpha,\infty)}(\phi)
i c_lpha \phi' ext{ in } \mathbb{R}, \hspace{0.3cm} \phi(0)=lpha, \hspace{0.3cm} \phi'(0)=0.$$

Claim: $\exists h_1 \in \mathbb{R}$ such that $\phi(x) = \phi_{\alpha}(x - h_1) \ \forall x \in \mathbb{R}$.

Thus there exists a sequence $t_n \to +\infty$ such that

$$\sup_{y\leq R} |u(y+c_lpha t_n,t_n)-\phi_lpha(y-h_1)| o 0 \quad ext{for} \ \ R>0.$$

Step 5 (Quasi-convergence global in space)

Recall U_{γ} and ϕ_{lpha} and use comparison argument to

 $|u(y+c_{\alpha}t_n,t_n)-\phi_{\alpha}(y-h_1)| \leq a_+-\min\{u(y+c_{\alpha}t_n,t_n),\phi_{\alpha}(y-h_1)\}.$

Final Phase Modifying the argument in [X. Chen '92], we shall prove exponential convergence of $u(\cdot, t)$ to $\phi_{\alpha}(\cdot - c_{\alpha}t)$ over \mathbb{R} as $t \to +\infty$ (without taking any subsequence). Set

$$w^{\pm}(x,t) := \phi_{lpha}(x - c_{lpha}t \pm \sigma\delta(1 - \mathrm{e}^{-eta t}) - h^{\pm}) \pm \delta\mathrm{e}^{-eta t}$$

for $h^{\pm} \in \mathbb{R}$, $\delta \in (0, \delta_0)$ and $\beta, \sigma > 0$. Then w^{\pm} turns out to be a superand a subsolution to $(AC)_{\alpha}$, provided δ_0, β, σ are small enough.

$\swarrow \text{Lemma 7 (Enclosing)} \\ \text{Set } h^+ = 0 \text{ and assume } \phi_{\alpha}(\cdot - h^-) - \delta \leq u(\cdot, 0) \leq \phi_{\alpha}(\cdot) + \delta \text{ in } \mathbb{R}. \\ \text{If } \delta \in (0, \delta_0) \text{ and } h^- \geq 0 \text{ are small enough, one can take } \varepsilon \in (0, 1) \\ \text{and } t \gg 1 \text{ such that} \end{cases}$

$$\phi_{lpha}(x-c_{lpha}t-arepsilon(h^{-}-\delta))-arepsilon\delta\leq u(x,t)\leq \phi_{lpha}(x-c_{lpha}t+arepsilon\delta)+arepsilon\delta.$$



By comparison principle, $w^-(x,t) \leq u(x,t) \leq w^+(x,t).$



Set $W^{\pm} := \pm (w^{\pm} - u) \geq 0$. Then

$$\partial_t W^\pm - \partial_x^2 W^\pm \geq \mp ig(f(w^\pm) - f(u)ig) \geq -M W^\pm.$$

By strong maximum principle, one has $W^+ \geq {}^\exists c_0 > 0$ or $W^- \geq c_0 > 0$.



On the intermediate region, $W^+ \geq c_0 > 0$ implies

$$u(x,t) \leq \phi_lpha(x-c_lpha t+\sigma\delta(1-\mathrm{e}^{-eta t})-\Delta)+\delta\mathrm{e}^{-eta t}$$

for some $\Delta > 0$ (\leftarrow improvement !).



If $\phi'_{lpha} \ll 1$, one can then enclose u(x,t) with smaller error:

$$\phi_{\alpha}(x-c_{\alpha}t) \approx \phi_{\alpha}(x-c_{\alpha}t-\Delta) - \Delta \phi_{\alpha}'(x-c_{\alpha}t)$$

On the left region, $0 < r_0 < h^-$ should be small enough.

Final remark

For sufficiently regular solutions, one can derive a motion equation for the free boundary r(t) as follows:

$$rac{\mathrm{d}r}{\mathrm{d}t}(t) = -rac{\partial_x^3 u(r(t),t)}{f(lpha)} \quad ext{ for } t>0.$$

cf.) Stefan problem

$$rac{\mathrm{d}r}{\mathrm{d}t}(t) = -\mu \partial_x u(r(t),t) \quad ext{ for } t>0$$

[Du-Lin'10][Du-Guo'11,'12][Du-Lou-Zhou'15][Du-Matsuzawa-Zhou'15][Kaneko-Yamada'11,'18][Kaneko-Matsuzawa'15,'18][Kaneko-Matsuzawa-Yamada'20]... cf.) The solution u(x, t) turns out to satisfy

$$u(r(t),t)=lpha, \;\; \partial_x u(r(t),t)=0, \;\; \partial_x^2 u(r(t)+0,t)=f(lpha)$$
 for any $t>0.$

Thank you for your attention !

Goro Akagi Tohoku University, JP

goro.akagi@tohoku.ac.jp